Constraints on the neutron star equation of state

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Motivation

- Stellar matter EoS: which NS properties do we get from a constrained EoS?
- How important is it to take a crust-core unified EoS?
- Hyperons in hot dense matter: what do the constraints tell us for equation of state?

Which quantities are required from consistent EoS by WG2?

Summary

- Constrained EoS
 - Experimental, theoretical and observational constraints

- Constrained EoS: consequences on the NS properties
- Constrained hyperonic EoS
 - Propoerties of hot dense matter with hyperons

Constraints on the EoS

- laboratory measurements of nuclear properties and reactions
 - nuclei, hypernuclei, heavy-ion collisions (HIC)
 - nuclei probe saturation and/or subsaturation densities of symmetric or almost symmetric nuclear matter
 - high density EoS from HIC: depends on transport models
- theoretical ab-initio calculations
 - neutron matter calculation from Quantum Monte Carlo and Chiral effective field theory up to saturation density due to perturbative behavior
- astrophysical observations
 - ► 2M_☉, R (still not well constrained), tidal deformability (pressure and energy density)
 - ► neutron star cooling and rotation (superfluidity)

Equation of state

Energy per nucleon

$$\boldsymbol{e}(\boldsymbol{\rho},\boldsymbol{\delta}) = \boldsymbol{e}(\boldsymbol{\rho},\boldsymbol{0}) + \boldsymbol{S}(\boldsymbol{\rho})\boldsymbol{\delta}^2$$

$$\rho = \rho_n + \rho_p, \, \delta = (\rho_n - \rho_p)/\rho$$

EoS for symmetric matter

$$e(\rho, 0) = e(\rho_0) + \frac{K}{2}x^2 + \frac{Q}{6}x^3 + \mathcal{O}(4), \quad x = \frac{\rho - \rho_0}{3\rho_0}$$

Density dependent symmetry energy

$$S(\rho)=J+Lx+\frac{K_{sym}}{2}x^2+\mathcal{O}(3),$$

- Properties of nuclear matter characterized by expansion coefficients
 - ▶ $\rho_0, B = m e(\rho_0), K, Q, M = Q + 12K, J, L, K_{sym}$

Nuclear constraints : Terrestrial experiments



• $B \sim 16$ MeV, $\rho_0 \sim 0.15 - 0.16$ fm⁻³

(nuclear masses, density distributions)

- $K = 230 \pm 40$ MeV (from analysis of ISGMR Khan PRL109)
 - ▶ but 250 < K < 315 MeV in Stone 2014
- constraints in J L plane:
 - ► J = 29.0 32.7 MeV, L = 44 66 MeV (Lattimer et al 2013,2014)
 - neutron skin thickness and giant dipole resonance of ²⁰⁸Pb, heavy ion collisions, measured nuclear masses, isobaric analog states, electric dipole polarizability α_D

► $J = 31.7 \pm 3.2$ MeV, $L = 58.7 \pm 28.1$ MeV (Oertel et al 2017)

Nuclear constraints



Neutron matter microscopic calculations

- chiral effective field theory constrain the properties of neutron matter up to ρ₀ (Hebeler et al 2010, 2013)
- realistic two- and three-nucleon interactions using quantum Monte Carlo techniques (Gandolfi et al 2012)

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Constraining the EOS from neutron star masses

massive neutron stars



Oertel et al arxiv:1610.03361

- ► PSR J0348+0432 (2.01 (4)M_☉, (advance of periastron) (Antoniadis et al, 2013)
- ▶ PSR J1614 $-2230~(1.928(17)M_{\odot},~(Shapiro delay)~(Demorest et all Compared on the second second$

2010, Fonseca et al. 2016)

► PSR J1946+3417 (1.828(22) M_☉, (Shapiro delay and advance of periastron) (Barr et al, MNRAS 2017)

Imposing 2*M*_☉ Fortin et al PRC 94,035804



- ► All EoS are causal and predict M > 2.M_☉
 - range of radii spanned:3km (1M_☉) and 4km (2M_☉)
- imposing lab and theoretical constraints:only 4 models remain
 - range of radii spanned:1km (1M_☉) and 2km (2M_☉)
 - large high mass uncertainty: lack of constraints on high density EoS!

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Nuclear constraints: $\Delta R_{1.4}$ and $\Delta \ell_{cr}$



(Fortin PRC94,035804(2016))

Imposing nuclear (experimental, theoretical, observational) constraints:

- $\Delta R_{1.0} = 0.9$ km, $\Delta R_{1.4} = 1.3$ km, $\Delta R_{2.0} = 2.3$ km
- ▶ R uncertainties reduced to 30% ($M_{1.0}$) and 50% ($M \ge 1.4M_{\odot}$)

Non-unified EOS

Radius and crust thickness



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(Fortin PRC94,035804(2016))

- core EOS: GM1, crust EOS: Douchin & Haensel
- different crust-core matching procedures
- ► 1.4 M_{\odot} star: ΔR =420 m, $\Delta \ell_{cr}$ = 350 m

Constraining the EOS from NS

Alam et al PRC 94, 052801



- correlation with 18 RMF EoS+ 24 Skyrme EoS, unified EoS, all describe 2M_☉ stars
- $P = \frac{\rho_0 x^2}{3} \left[K_0(x-1) \left(1 \frac{2x}{3} \right) + \frac{M_0}{18} (x-1)^2 + L_0 \delta^2 \right] .,$ $M_0 = Q_0 + 12K_0, \qquad x = \rho/\rho_0$
- $M_0(n_0) = 1800 2400 \text{ MeV}$ from energies ISMGR (De 2015)
- Prediction: R_{1.4} = 11.09 12.86 km

Constraining the EOS at high densities



HIC: there are still no consensus in data and transport codes analysis

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- ► GW170817: constraints on tidal deformability $\Lambda = \frac{2}{3}k_2 \left(\frac{R}{M}\right)^5$ $\Lambda < 800(500)$ at 90% (50%) confidence
 - Λ_1 : high mass M_1 satisfies 1.365 < M_1 < 1.60 M_{\odot}
 - Λ_2 : low mass M_2 from $M_{\rm chirp} = 1.188 M_{\odot}$.

Entrainment matrix for β equilibrium matter

Effect of the symmetry energy slope L and of the magnetic field B

Superfluid currents at T = 0: Q_i momentum per particle of Cooper pairs i (Gusakov, Kantor, Haensel PRC79)

$$\begin{aligned} \boldsymbol{j}_n &= \boldsymbol{Y}_{nn} \boldsymbol{Q}_n + \boldsymbol{Y}_{np} \boldsymbol{Q}_p \\ \boldsymbol{j}_p &= \boldsymbol{Y}_{pn} \boldsymbol{Q}_n + \boldsymbol{Y}_{pp} \boldsymbol{Q}_p. \end{aligned}$$



Adiabatic index for β equilibrium matter

Effect of the symmetry energy slope L and of the magnetic field B



Effect of the magnetic field B on the crust thickness



diffuse separation between homogeneous and clusterized matter

A D > A P > A D > A D >

• transition thickness $\Delta \rho > 0.01$ fm⁻³ for $B = 2.2 \times 10^{15}$ G

Fang PRC95.045802

Effect of the magnetic field B on the crust thickness

Crust thickness and symmetry energy



 $T = 0, B = 5 \times 10^{16} \text{G}$

- Transition densities
- ▶ crust thickness ∆R
- crust fractional I
- ▶ B = 0 (black)

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Magnetized Crust thickness versus Temperature



NL3 $\omega \rho$, L = 55 MeV

- Transition densities
- ▶ crust thickness ∆R
- crust fractional I

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How important are hyperons in NS?



- ► Is it possible to satisfy hypernuclei properties and have $2M_{\odot}$ NS?
- What do experimentally constrained hyperonic EoS tell us?

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Hypernuclei



- ► scattering events: → not enough to constrain interactions
- hypernuclei
 - \gtrsim 40 single Λ -hypernuclei
 - a few double Λ and single- Ξ
 - ► no unambiguous ∑-hypernucleus:most probably ∑-nucleus potential repulsive

• attractive Ξ -nucleus interactions, $U_{\Xi}^{N} \sim -14$ MeV.

AA binding from double and single A-hypernuclei

$$\Delta B_{\Lambda\Lambda} = B_{\Lambda\Lambda} \begin{pmatrix} A \\ \Lambda\Lambda \end{pmatrix} - 2 B_{\Lambda} \begin{pmatrix} A^{-1} \\ \Lambda \end{pmatrix}$$

Unambiguous measurement ⁶_{AA}He by KEK (2001)

$$\Delta B_{\Lambda\Lambda}$$
 = 0.67 \pm 0.17MeV

Hyperonic stars

Fortin et al PRC 95



Vector meson couplings

choice a: SU(6) symmetry for ω , varying ϕ -hyperon **choice b**: $g_{Y\omega} = g_{N\omega}$, varying ϕ -hyperon ρ -meson: $g_{\rho\Xi} = \frac{1}{2}g_{\rho\Sigma} = g_{\rho N}$

• Σ - σ coupling: $U_{\Sigma}^{N}(n_{0}) = 0, +30 \text{ MeV}$

► Ξ - σ coupling: $U_{\Xi}^{N}(n_{0}) = -14$ MeV and $U_{\Xi}^{N}(2n_{0}/3) = -14$ MeV

AN-potential

Fortin et al PRC 95

Model	$R_{\omega\Lambda}$	$R_{\sigma\Lambda}$	$U_{\Lambda}^{N}(n_{0})$
TM1-a	2/3	0.621	-30
TM1-b	1	0.892	-31

DDME2D-a	2/3	0.621	-32
DDME2D-b	1	0.896	-35

YN potential

$$U_Y^N(n_0) = - \left(g_{\sigma Y} + g'_{\sigma Y} \rho_s\right) \sigma_0 + \left(g_{\omega Y} + g'_{\omega Y} n_0\right) \omega_0,$$

- ► U^N_A(n₀) ≃ -30 MeV in agreement with the binding energy of single A-hypernuclei in the s- and p-shells
- *R*_{σΛ} ~ 0.62 for SU(6) value for g_{ωΛ}, independent of the model considered.

$\Lambda\Lambda$ -potential Fortin et al PRC 95

Model			$\Delta B_{AA} = 0.50$		$\Delta B_{\Lambda\Lambda} = 0.84$	
		$oldsymbol{R}_{\phi \Lambda}$	$R_{\sigma^*\Lambda}$	$U^{\Lambda}_{\Lambda}(n_0)$	$R_{\sigma^*\Lambda}$	$U^{\Lambda}_{\Lambda}(n_0)$
TM1-a	SU(6)	$-\sqrt{2}/3$	0.533	-11.2	0.557	-14.2
TM1-b		$-\sqrt{2}/2$	0.843	2.7	0.864	-1.2
NL3-a	SU(6)	$-\sqrt{2}/3$	0.534	-9.9	0.559	-13.2
NL3-b		$-\sqrt{2/2}$	0.846	9.0	0.868	4.8
DDME2D-a	SU(6)	$-\sqrt{2}/3$	0.535	-11.9	0.555	-11.7
DDME2D-b		$-\sqrt{2}/2$	0.846	-3.4	0.862	-3.4

A potential in pure A matter :-14 < $U_A^A(n_0)$ < +9 MeV in literature taken between -1 or -5 MeV

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Protoneutron stars/Binary neutron star mergers

- PNS and binary NS mergers:
 - High densities and high temperatures are attained inside
 - ► T = 10 100 MeV affect the NS composition, favor the production of non-nucleonic degrees of freedom
 - Which is the role of non-nucleonic degrees of freedom?
- How will constraints on the EoS affect the evolution of PNS and NS mergers in the evolution of these systems?
- Information on the EoS is expected from
 - the tidal deformability during late inspiral
 - post merger oscillations
 - observation of GW correlated with an electromagnetic signal

Hyperons in hot dense matter

What do the constraints tell us for the EoS?

Symmetry enegy



(Oertel et al 2017)

contraint: ab-initio calculations of pure neutron matter



Hyperon fractions $T = 30 \text{ MeV}, Y_Q = 0.1, 0.3, 0.5$



Temperature for fixed entropy per baryon $s_B = 1, 2, 4$



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Hyperons in hot stellar matter: conclusions

- DD2Y EoS (Marques2017) and SFHoY: only general purpose EoS models with the entire baryon octet and compatible with the relevant constraints on the EoS
- ► SFHo:softer symmetry energy than DD2, additional repulsion
 - SFHoY has smaller hyperon fractions
 - the effects on thermodynamic properties are much less pronounced
- Possible consequences
 - different proto-neutron star evolution
 - different impact of hyperonic degrees of freedom on neutron star merger dynamics

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► SFHoY and SFHoY*: available in COMPOSE database

Thank you !

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